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# The *HST*/ACS Coma Cluster Survey – X. Nuclear star clusters in low-mass early-type galaxies: scaling relations

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## ABSTRACT

We present scaling relations between structural properties of nuclear star clusters and their host galaxies for a sample of early-type dwarf galaxies observed as part of the *Hubble Space Telescope* (*HST*) Advanced Camera for Surveys (ACS) Coma Cluster Survey. We have analysed the light profiles of 200 early-type dwarf galaxies in the magnitude range  $16.0 < m_{F814W} < 22.6$  mag, corresponding to  $-19.0 < M_{F814W} < -12.4$  mag. Nuclear star clusters are detected in 80 per cent of the galaxies, thus doubling the sample of *HST*-observed early-type dwarf galaxies with nuclear star clusters. We confirm that the nuclear star cluster detection fraction decreases strongly towards faint magnitudes. The luminosities of nuclear star clusters do not scale linearly with host galaxy luminosity. A linear fit yields  $L_{\text{nuc}} \sim L_{\text{gal}}^{0.57 \pm 0.05}$ . The nuclear star cluster–host galaxy luminosity scaling relation for low-mass early-type dwarf galaxies is consistent with formation by globular cluster (GC) accretion. We find that at similar luminosities, galaxies with higher Sérsic indices have slightly more luminous nuclear star clusters. Rounder galaxies have on average more luminous clusters. Some of the nuclear star clusters are resolved, despite the distance of Coma. We argue that the relation between nuclear star cluster mass and size is consistent with both formation by GC accretion and *in situ* formation. Our data are consistent with GC inspiralling being the dominant mechanism at low masses, although the observed trend with Sérsic index suggests that *in situ* star formation is an important second-order effect.

**Key words:** galaxies: dwarf – galaxies: nuclei – galaxies: star clusters: general.

## 1 INTRODUCTION

Nuclear star clusters (NSCs) are dense concentrations of stars in galaxy centres. Although their sizes, as measured in nearby spiral

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galaxies, are found to be comparable to globular clusters (GCs; Böker et al. 2004), their luminosities can exceed the brightest Milky Way GC by orders of magnitude. NSCs have been found in galaxies of all Hubble types, from sub- $L_*$  early-type galaxies (Côté et al. 2006) to late-type spiral galaxies (Carollo et al. 1997; Carollo, Stiavelli & Mack 1998; Carollo 1999; Böker, van der Marel & Vacca 1999; Matthews et al. 1999; Böker et al. 2001, 2002; Carollo et al. 2002; Georgiev & Böker 2014), early-type spiral bulges (Balcells et al. 2003; Balcells, Graham & Peletier 2007) and in early-type dwarf galaxies (dEs;<sup>1</sup> Graham & Guzmán 2003; Lotz, Miller & Ferguson 2004; Grant, Kuipers & Phillipps 2005; Côté et al. 2006). Also the Milky Way is known to have an NSC (Launhardt, Zylka & Mezger 2002; Schödel et al. 2007; Schödel, Merritt & Eckart 2009).

There are several formation scenarios for NSCs, which can essentially be divided into two main scenarios: build-up through the accretion of (globular) star clusters, or by star formation *in situ*.

The GC accretion scenario was first developed by Tremaine, Ostriker & Spitzer (1975). In this scenario, GCs gradually approach the galaxy centre as angular momentum is slowly removed by dynamical friction with stars. At the centre, the GCs merge to form an NSC. This model was later refined to account for the long dynamical friction time-scales (e.g. Milosavljević 2004). Because of their low-velocity dispersion, dynamical friction of star clusters should be most efficient in low-mass dwarf galaxies (e.g. Lotz et al. 2001).

In the *in situ* formation scenario, gas is transported to the centre of a galaxy, where it cools and forms an NSC. Several mechanisms to transport gas to the inner parsecs of gas-rich galaxies have been proposed (e.g. Mihos & Hernquist 1994; Milosavljević 2004; Hopkins & Quataert 2011). Since simulations and calculations involving hydrodynamics are more complex than pure gravitational computations, the theoretical framework for the dissipational formation of NSCs provides less quantitative predictions than the inspiralling cluster scenario.

There is evidence that both the *in situ* as well as the GC inspiralling formation happen in galaxies. Evidence for the GC inspiral scenario is seen in the lack of GCs in the inner parts of dE galaxies (Lotz et al. 2004). However, the colours of some of the NSCs studied by these authors are inconsistent with the colours of the old populations typically seen in GCs. Similarly, NSCs in spiral galaxies show often evidence for recent star formation (Walcher et al. 2005). Evidently, there is no consensus on the dominant formation channel of NSCs in dEs. And if such a dominant formation channel exists, it may well be dependent on the morphological type or mass of the host galaxy or indirectly on the environment.

For the spiral galaxy NGC 4244, Hartmann et al. (2011) infer from a detailed dynamical study that at least 50 per cent of the mass of the NSC was formed by gas accretion, but also that at least 10 per cent of the stars have to have formed outside the cluster. Such detailed studies are currently only possible for nearby galaxies. However, at larger distances, the study of scaling relations of NSCs may provide insights into the formation mechanisms.

The luminosities of NSCs in dEs are known to scale with the host galaxy luminosity (Graham & Guzmán 2003). This relation was also analysed by Grant et al. (2005), Côté et al. (2006), Turner et al. (2012) and Scott & Graham (2013), although, due to sample selection, only for bright dEs ( $M_B \lesssim -15$ ). The same groups have studied the sizes of NSCs in dEs as a function of NSC luminosity.

It is unclear what these scaling relations look like for NSCs in low-mass dEs. In this paper, we address the scaling relations of NSC in low-mass early-type galaxies through the analysis of the light distribution of 200 dEs in the Coma cluster and try to infer the dominant formation mechanism for the NSCs. A majority of these dEs are located (in projection) close to the core of the cluster.

The GC inspiral scenario provides testable predictions for the mass-ratio of NSCs and host galaxies, as well as for the geometrical sizes of the NSCs, although not without some uncertainty, since also a widely accepted theory for the formation of GCs is still lacking. Bekki et al. (2004) present scaling relations for the sizes of NSCs based on  $N$ -body simulations of the merging of equal-mass star clusters; however, they do not predict a relation between host galaxy mass and NSC mass. Antonini (2013) on the other hand contains a comprehensive overview of scaling relations for both formation mechanisms, most of which are analytically derived.

One of the few qualitative predictions for the relation between NSC mass and host galaxy velocity dispersion for *in situ* formation is given by McLaughlin, King & Nayakshin (2006), where it is assumed that the winds from giant stars and supernovae in the star cluster provide feedback to the galaxy, in a way that is similar to supermassive black hole (SMBH) feedback. There is however direct evidence for episodic gas accretion into NSCs in spiral galaxies. Emsellem & van de Ven (2008) analyse the tidal forces in Sérsic galaxies, and find that in galaxy centres on the scale of the size of typical NSCs, tidal forces become disruptive in galaxies that are too concentrated (i.e. Sérsic index  $n > 3.5$ ). They derive that the mass of a central massive object (CMO) that would remove all compressive tidal forces scales linearly with host galaxy mass. Although reaching this maximum CMO mass depends on the amount of gas available for star formation, if indeed tidal forces are a key ingredient in the formation of NSCs, the formation efficiency should depend negatively on the Sérsic index.

It has been suggested that NSCs and central SMBHs follow a  $M_{\text{cmo}}-\sigma$  scaling relation with the same slope (Ferrarese et al. 2006; Wehner & Harris 2006), although this has been challenged (see for example, Balcells et al. 2007; Erwin & Gadotti 2012; Graham 2012; Leigh, Böker & Knigge 2012; Neumayer & Walcher 2012; Scott & Graham 2013). The existence of the  $M_{\text{cmo}}-\sigma$  relation implies that the evolution of NSCs and SMBHs may possibly somehow be linked. Likely there is some interaction between NSCs and SMBHs, such that one may prevent the growth of the other or even destroy it (McLaughlin, King & Nayakshin 2006; Merritt 2009; Nayakshin, Wilkinson & King 2009). The Milky Way was the first galaxy known to host both an NSC and an SMBH. Although this could be interpreted as a peculiar case, it has been shown that also in other galaxies NSCs and SMBHs coexist (Seth et al. 2008; Graham & Spitler 2009).

Very massive ellipticals have been found to lack NSCs. Merritt (2009) provides a mechanism to dissolve NSCs by absorption of energy from the surrounding galaxy, and shows that the addition of an SMBH to the NSC always leads to expansion of the NSC. Bekki & Graham (2010) suggest that a merger of intermediate-mass elliptical galaxies in which NSCs and black holes (BHs) coexist will result in an elliptical galaxy in which the newly formed star cluster is structurally and dynamically altered so that it may easily get destroyed. However, Antonini (2013) provides a different explanation for the lack of NSCs in galaxies with SMBHs (the argument was derived for dissipationless formation): the tidal forces from the SMBH rip inspiralling clusters apart at large radii from the centre, so that the NSC ends up having a much lower density than in the absence of an SMBH and may therefore not be properly disentangled from the host galaxy.

<sup>1</sup> We use the term dE both for dwarf elliptical as well as dwarf lenticular galaxies.

Whatever physical interaction is going on between NSCs and SMBHs, both trace the mass accretion to the central parsecs of a galaxy, and the study of NSCs can therefore provide insights in the formation of SMBHs and their host galaxies, as well as their haloes.

Due to its exquisite resolution, the *Hubble Space Telescope* (*HST*) is a particularly good instrument for studying NSCs. The work of Lotz et al. (2004) and Côté et al. (2006) has shown that ground-based observations of nuclear clusters are often of too low resolution for studying remote NSCs, because a significant fraction of the clusters is missed. So far, only a small set of NSCs in dEs in Coma has been studied (Graham & Guzmán 2003). Coma is the *most convenient* very rich (Abell class 2) and dense cluster to study in the local Universe. In the centre, the projected density of faint galaxies is almost a magnitude higher than in Virgo (Weinmann et al. 2011). If the environment plays a role in shaping NSCs, Coma is the place where this could be measured best.

For this paper, we focus on the scaling relations of a sample of dwarf early-type galaxies in the Coma Advanced Camera for Surveys (ACS) survey. Despite the distance to Coma, we manage to double the number of known NSCs in dEs observed with *HST*.

In Sections 2 and 3, we briefly summarize the Coma ACS survey, data reduction and analysis. Section 4 presents some basic results of this analysis. We discuss these results in Section 5.

## 2 OBSERVATIONS AND DATA REDUCTION

The Coma ACS Survey (Carter et al. 2008, henceforth C08) provides data in two passbands for 25 fields pointed at the core of the Coma cluster and at the outskirts. The exposure times in the two passbands,  $F814W$  and  $F475W$  (which are roughly equivalent to the  $I_C$  and  $g$  band) were  $\sim 1400$  and  $\sim 2600$  s. The original envisaged coverage of the cluster was much larger than 25 fields, but due to the ACS failure in 2008 January the survey was not completed. Of the 25 observed fields, five fields have data of lower sensitivity (visits 1, 3, 12, 13 and 14, see also C08) because of a lack of data or poor quality.

The pixel scale of the ACS data reduced with the drizzling scheme as described in C08 is 0.05 arcsec per pixel. Throughout this paper, we assume a distance to Coma of 100 Mpc (see C08). This corresponds to  $\sim 25$  pc per pixel at the distance of Coma. The  $10\sigma$  detection limit of point sources is 26.8 in  $F814W(AB)$  and 27.6 in  $F475W(AB)$ , although this is in optimal conditions.

### 2.1 Sample

An advantage of observing the Coma cluster compared to nearby clusters such as the Virgo cluster, is that the relative distance uncertainty between galaxies in the cluster is low. The disadvantage is that membership of galaxies is not always well determined.

In Hammer et al. (2010) more than 70 000 sources were detected in the  $F814W$  passband along the line of sight of Coma. Most of these sources are either GCs in Coma or background galaxies. The best way to establish cluster membership is by spectroscopically confirming that a galaxy has a redshift consistent with the velocities of other cluster members. However, spectroscopic observations have targeted mainly bright galaxies. Luckily, faint dE cluster members stand out against the background because of their low surface brightness.

We use a subset of the sample defined in Trentham et al. (in preparation, Paper XI), which is used for determining the luminosity function of galaxies in Coma. The selection of this sample was based on by-eye identification of galaxies on the ACS frames. The catalogue is complete, except at faint magnitudes. Faint galax-

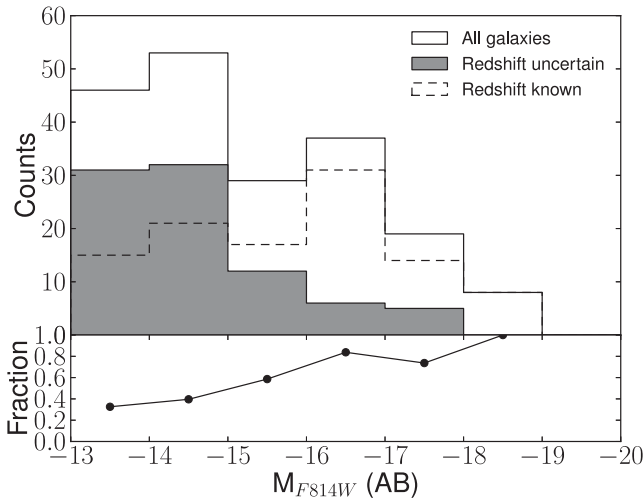
ies ( $M_{F814W} = -10$ ) with similar structural properties as the Local Group galaxy Draco should be detected in the ACS frames. However, we exclude very faint dEs to avoid the introduction of strong biases in the sample of nuclear clusters: it has been found that nuclear clusters are generally 5–6 mag fainter than the host galaxy (see also Section 4.3) so that detection of NSCs is probably not possible in the faintest galaxies, whether or not these galaxies are really nucleated.

At faint magnitudes, the dominant noise is a combination of the readnoise and the sky background noise with a small contribution of Poisson noise from the host galaxy (see also Section 4.3 and Fig. 6). For a host galaxy absolute magnitude  $M_{F814W} = -13$  mag, the difference between the host galaxy magnitude and the magnitude of a point source that could be detected at  $3\sigma$  level is expected to be  $\sim 4.7$  mag (and lower than 4 for  $M_{F814W} = -11$ ).

We decided therefore to clip the faint part of our sample at  $m_{F814W}(AB) = 22.6$  ( $M_{F814W}(AB) = -12.6$ ,  $M_B \approx -11.0$ ), also because obtaining reliable structural parameter fits of the host galaxy becomes difficult for galaxies that are fainter than this magnitude. We use the  $F814W$  magnitudes for clipping since it is a better proxy of stellar mass than the  $F475W$  magnitudes. We also introduce a limiting magnitude on the bright side of the magnitude range since for this paper we focus exclusively on NSCs in dE galaxies: it has been found that low-mass (non-dwarf) ellipticals often have a central light excess, but lack the distinct surface brightness bump and often distinct colours seen in lower luminosity galaxies (e.g. Côté et al. 2006; Kormendy et al. 2009). Although scaling relations of the luminosity and size of this nuclear excess light and NSCs with host galaxy properties may be continuous, the formation mechanisms are not necessarily the same. The other reason for discarding high-mass galaxies is that they often require more than two components for the fitting (e.g. Aguerri et al. 2004; Dullo & Graham 2014; Huang et al. 2013; Head et al. 2014; Janz et al. 2014; Läsker, Ferrarese & van de Ven 2014; Weinzirl et al. 2014). We therefore only use galaxies fainter than  $M_{F814W}(AB) = -19.0$ .

Our sample contains both spectroscopically confirmed members and possible members, which were identified by eye (by NT and HF). In Paper XI, these possible members are categorized into three classes: almost certainly members (class 1), likely members (class 2) and possible members (class 3). Spectroscopic follow-up of a subset of possible members in the Paper XI sample has shown that the contamination of non-members in class 3 is around 50 per cent, and approximately 10 per cent for sources in class 2 (Chiboucas et al. 2010). A distribution of confirmed and non-confirmed sources over magnitude can be found in Fig. 1. The number of sources in our magnitude-limited sample with successful structural parameter fits that are in class 0 (confirmed members) is 104, in class 1 is 13, in class 2 is 58 and in class 3 is 23. In this paper, we analyse galaxies in all classes, but since membership is uncertain for class 3 objects, we exclude them from our measurements of scaling relations. For clarification, we note that whenever we refer to the ‘full sample’, we mean the magnitude-limited sample containing objects from all four classes, whereas with ‘likely members’ and ‘probable members’ we denote the magnitude-limited sample of galaxies in classes 2 and lower. We note that a number of galaxies were excluded from our analysis, because they either had late-type or irregular features such as interstellar medium or significant spiral arms, were compact ellipticals, or were very difficult to fit, usually because they had a close bright companion. These sources are listed in Appendix A. Similarly, due to the chosen dithering scheme, galaxies close to the edges of the frame suffer from lower signal-to-noise ratio (S/N) and cosmic rays. A handful of these galaxies also ended up in Appendix A.





**Figure 1.** Counts of galaxies with known and unknown redshift the fraction of those as a function of absolute magnitude of the host galaxy (for likely member galaxies). The black histogram shows the distribution of sources without redshift information as a function of host galaxy magnitude. The dashed histogram shows the sources which are confirmed member galaxies.

For sources that appeared in multiple frames, we chose to fit the source in the frame with the best S/N. We note that, primarily due to the curtailment of the ACS survey, 90 per cent of our sample galaxies are at a projected distance of 500 kpc from the core, whereas the virial radius of Coma is about 3 Mpc (Łokas & Mamon 2003).

### 3 ANALYSIS

#### 3.1 Choice and justification of surface brightness profile parametrization

The stellar populations of NSCs are distinct from those of the host galaxy (e.g. Lotz et al. 2004), so that treating the host galaxy and the NSC as distinct morphological components is justified. We use the Sérsic model (Sersic 1968) to describe the light profile of the host galaxy. This profile is a generalization of the exponential profile, which has been used to fit disc galaxies (Patterson 1940; Freeman 1970), and the  $R^{1/4}$ -profile, commonly used for elliptical galaxies (de Vaucouleurs 1948). The Sérsic profile is given by

$$I(R) = I_e \exp \left\{ -b_n \left[ (R/R_e)^{1/n} - 1 \right] \right\}$$

and provides an accurate description of the light profile of low-mass early-type galaxies outside the centre (Caon, Capaccioli & D’Onofrio 1993; D’Onofrio, Capaccioli & Caon 1994).  $I_e$  is the surface brightness at the effective radius. The Sérsic index  $n$  defines the curvature of the profile; profiles with high  $n$  are more cuspy.  $b_n$  plays the role of a normalization constant, and can be approximated by  $b_n \approx 0.19992n - 0.3271$  for  $0.5 < n < 10$  (Capaccioli 1989). A comprehensive review of the Sérsic model can be found in Graham & Driver (2005).

If NSCs in Coma follow the same size distribution as those in Virgo (Côté et al. 2006), we expect most NSCs to be unresolved. Our first choice of modelling is therefore to fit them as a point spread function (PSF) convolved point source. However, since some of them may be marginally resolved, we also fit them with a Gaussian, similar to Graham & Guzmán (2003) and Kourkchi et al. (2012). Although there is no physical justification for using a Gaussian profile, the use of a more complicated profile is not warranted by

our data, and the Gaussian profile has one fewer free parameters than a King model.

#### 3.2 Fitting code

We make use of BAGATELLE (see Appendix B), a two-dimensional Bayesian fitting code for modelling surface brightness images of galaxies. The advantage of this code over other codes is that it allows for (i) a full exploration of the posterior distribution, which can help quantify any degeneracies, and (ii) a quantitative way of deciding which profile fits a galaxy best; given two model profiles with priors on their parameters, the code calculates and compares the marginalization over all variables, known as the Bayesian *evidence*, of the profiles. We describe the fitting code in more detail in Appendix B and give a comparison between the results of fits carried out by Weinzirl et al. (2014) with GALFIT and our fits in Appendix C.

#### 3.3 Masking

Neighbouring stars and galaxies were masked according to the SEXTRACTOR segmentation maps. As masking in this way was generally not sufficient, the SEXTRACTOR mask apertures were enlarged to an ellipse with major and minor axes of  $4 \times A\_IMAGE$  and  $4 \times B\_IMAGE$ . As close neighbours are sometimes missed by SEXTRACTOR, we manually checked the masks and tweaked them where necessary.

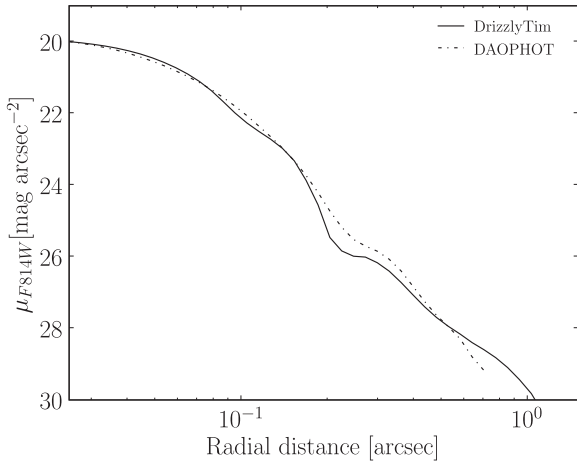
#### 3.4 Point spread function

When fitting sources which are unresolved or marginally resolved, detailed knowledge of the PSF is crucial. For our analysis, we have chosen to use the artificial PSFs from TINYTIM, drizzled in the same way as our HST/ACS data using a script called DRIZZLYTIM<sup>2</sup> (for details see Hoyos et al. 2011). Given the importance of the PSF, we have also determined empirical PSFs. The PSF of the ACS camera varies with position on the chip because the camera is mounted off-axis, and varies also as a function of time, since thermal fluctuations change the focus of the telescope (e.g. Rhodes et al. 2007).

To construct the empirical PSFs, we used archival data of the Galactic GC 47 Tuc. The exposure time of both the Coma and 47 Tuc observations are shorter than one HST orbit, which is the main time-scale of variation (Rhodes et al. 2007). Moreover, the time between the observations of the two targets is longer than six months. Although the empirical PSF will not be the optimal representation of the real PSF, it will at least show us how the uncertainty in the shape of the PSF affects our analysis.

The drizzled 47 Tuc data were processed using the DAOPHOT II package in IRAF (e.g. Stetson 1987). In overview and summary, we detected stars with the task daofind and calculated a PSF using 669 stars which were relatively bright, isolated and not saturated, with the task psf. After calculation of a PSF, stars near the PSF stars were subtracted from the image and the PSF was recalculated. We used a few such iterations, where during the last iterations we let the PSF vary quadratically as a function of position. We assumed a Moffat profile with  $\beta = 1.5$  for the analytic component of the PSF, although other analytic profiles fitted equally well. A comparison between radial surface brightness profiles of the empirical PSF and the DRIZZLYTIM PSF is given in Fig. 2. The empirical PSF is

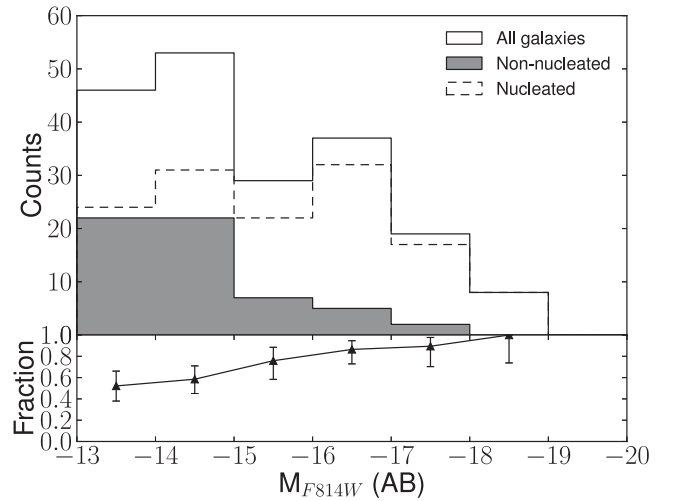
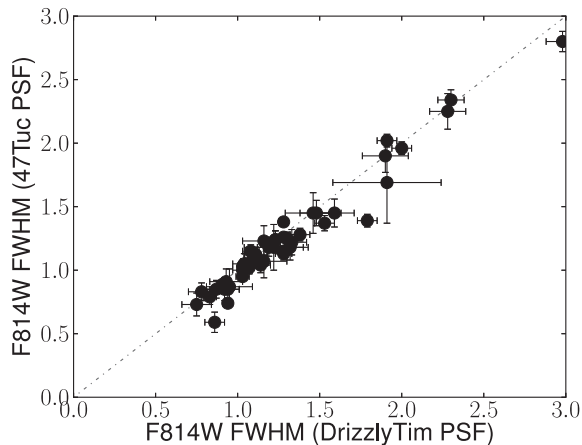
<sup>2</sup> DRIZZLYTIM is written by Luc Simard.



**Figure 2.** Normalized profiles of the  $F814W$  psf at the location  $x, y = (3300, 2400)$ . The solid line is the DRIZZLYTIM PSF, the dash-dotted line is the empirical PSF.

slightly broader than the theoretical PSF. The mismatch between the profiles of the two PSFs in the outer parts is a consequence of estimating the sky close to PSF stars because of crowding in the 47 Tuc observations. For detecting and measuring sizes of star clusters, and for the photometry of the main body of the galaxy, these wings are not relevant. However, it may cause us to underestimate the flux of an NSC. Based on the differences between the empirical and TINYTIM PSF in both  $F475W$  and  $F814W$ , we expect a systematic flux error of 2 per cent in magnitude, and less so in the colour, since it is a differential measurement.

During the course of this work, we found that several NSCs were resolved. To ensure that the sizes of the NSCs were not due to problems with the PSF, we measured sizes with both the empirical and theoretical PSFs. The left-hand panel of Fig. 3 shows that the measured sizes are not dependent on which PSF is used. To check that the sizes are not due to data-related issues, we show in the right-hand panel of Fig. 3 the sizes of resolved star clusters in the  $F814W$  band versus those in the  $F475W$  band. Although there is considerable scatter, the correlation suggests that the clusters are indeed resolved.

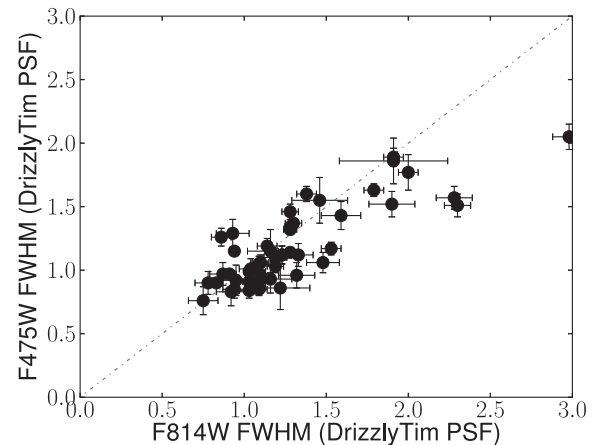


**Figure 4.** Counts of nucleated and non-nucleated galaxies and frequency of nucleation as a function of absolute magnitude of the host galaxy (for likely member galaxies). The black histogram shows the distribution of non-nucleated sources as a function of host galaxy magnitude. The dashed histogram shows the nucleated galaxies. Since faint NSCs may have been missed, the nucleation fraction is denoted as a lower limit for the points in the lower panel. The error bars denote the 95 per cent Jeffreys interval.

## 4 RESULTS

### 4.1 Detection fraction

Fig. 4 shows a histogram of galaxy counts as a function of absolute  $F814W$  magnitude, for nucleated and non-nucleated galaxies. We confirm the result of Graham & Guzmán (2003) and Côté et al. (2006) that the NSC fraction in bright dwarf galaxies is close to unity. Faint dEs show a lower fraction of nucleation, qualitatively consistent with results of Sandage, Binggeli & Tammann (1985) and van den Bergh (1986). However, it is not inconceivable that we have missed faint nuclei; the nucleation fractions in faint dEs are therefore likely lower limits. In fact, injecting the faintest genuine nucleated galaxy (VCC1895) from the ACS Virgo Cluster Survey (VCS) with updated structural parameters and noise added in a Coma image, we do not recover the nucleus. Over the full luminosity range, the fraction of non-nucleated galaxies is 16 per cent, similar



**Figure 3.** Left:  $F814W$  sizes of NSCs (in pixels) as measured with the empirical PSF versus  $F814W$  sizes of NSCs as determined with a theoretical PSF. Right:  $F475W$  sizes of NSCs (in pixels) versus  $F814W$  sizes of NSCs. Both sizes were measured using a DRIZZLYTIM PSF convolved with a Gaussian, while in parallel fitting the host galaxy with a Sérsic profile.

to the 13 per cent found by Graham & Guzmán (2003) for luminous dEs in Coma.

Even though this number is in agreement with the number found by Côté et al. (2006), we note that the average galaxy luminosity of the Virgo sample is higher. The samples have mean and median magnitude difference of 4.0 and 3.8 mag (2.8 and 2.7 for the nucleated galaxies in both samples). Given that our sample is volume and luminosity complete, it is representative of the luminosity function of dEs in the core of the Coma cluster.

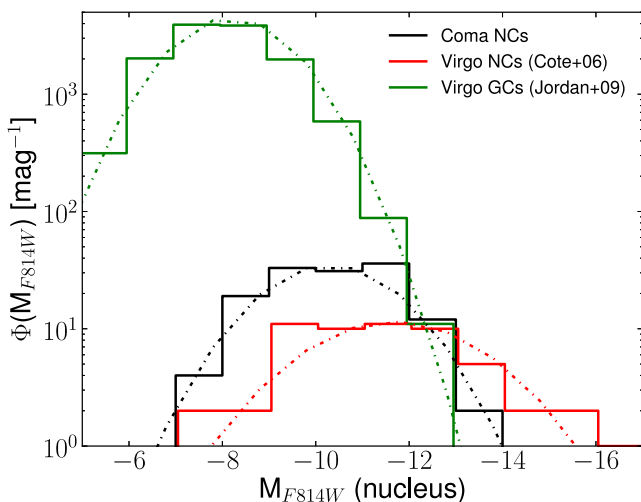
As a possible caveat, we note that the detection fraction of NSCs in visits 3 and 13, both with only half the planned exposure time, is lower than for the visits with complete exposure times in the core of the cluster. The total detection fraction of NSCs between  $-16 > M_{F814W} > -18 = 92 \pm 7$  per cent is consistent with the detection fraction of  $85 \pm 5$  per cent in the  $z$  band for galaxies in the same magnitude range in Côté et al. The small difference in detection fraction can probably be explained by different sample properties, since we have excluded galaxies with late-type features.

## 4.2 Luminosity function

In Fig. 5, we show the luminosity function of the NSCs, together with the luminosity function of NSCs in the Virgo cluster from Côté et al. and the luminosity function of GCs in Virgo (Jordán et al. 2009). The luminosity function is well described by a Gaussian, although we note that this functional form for the luminosity function is somewhat ad hoc and depends on survey depth (Turner et al. 2012).

The nuclei in our sample are slightly fainter than the NSCs in the ACS VCS sample (by about 1.5 mag in the  $F814W$  band), which is a consequence of the Coma sample containing more low-mass galaxies.

The faintest NSCs in our sample coincide with GCs at the turnover magnitude. For both formation scenarios there are qualitative arguments why there could be a lower limit to the NSC mass: in the accretion scenario, low-mass GCs have much longer dynamical friction time-scales, thus the formation of low-mass NSCs is suppressed; in the *in situ* scenario, infant mortality may be higher for low-mass star clusters in addition to feedback effects related



**Figure 5.** Luminosity function of NSCs in the  $F814W$  band. Also shown are the NSCs of galaxies in Virgo (Côté et al. 2006, red) and GCs in Virgo (Jordán et al. 2009, green), with the best-fitting Gaussians shown as dash-dotted lines.

to putative central BHs. The fact that the faintest NSCs coincide with the turnover magnitude of GCs in Virgo may also be coincidence. We note however that, except for the very brightest nuclei ( $F814W \approx -14$ ), most NSC magnitudes are consistent with being drawn from the massive end of the GC luminosity function.

## 4.3 Luminosity scaling relations

Several authors have pointed out a relationship between NSC luminosity and host galaxy luminosity in dEs (e.g. Graham & Guzmán 2003; Lotz et al. 2004; Grant et al. 2005; Côté et al. 2006). With our data set, we double the number of measurements for this relationship. Fig. 6 shows the NSC magnitude as a function of host galaxy magnitude in the  $F814W$  band. The host galaxy magnitude and NSC magnitude were determined as described in Section 3. The magnitude of the host galaxy was determined by extrapolating the Sérsic profile to infinity, and does not include the NSC. Because the NSC is on average 5–6 mag fainter, this should only make a marginal difference ( $\sim 0.01$  mag) in the magnitude of the host galaxy. The region where the detection limit for point sources is below  $10\sigma$  is indicated by the grey area below the dotted line. At low magnitudes, the sky dominates in the detection efficiency, whereas at high magnitudes, the Poisson noise of the centre of the host galaxy is dominant. In the same plot, we also show data from NSCs in the ACS VCS (Côté et al. 2006) and the ACS Fornax Cluster Survey (Turner et al. 2012). The conversion between  $F850LP$  and  $F814W$  is small (generally less than 0.2 mag) and holds for both host galaxy and star cluster – it is therefore unlikely that increased scatter is due to this conversion.

The magnitude–magnitude relation of Coma cluster NSCs and their host galaxies actually extends the relation found for Virgo and Fornax towards fainter magnitudes, although with a change of slope. Interestingly, the nucleus–galaxy luminosity ratio is higher than the one found by Côté et al. Performing a linear fit with fixed slope of one, we find (for the  $F814W$  band), for weighted and unweighted fits:

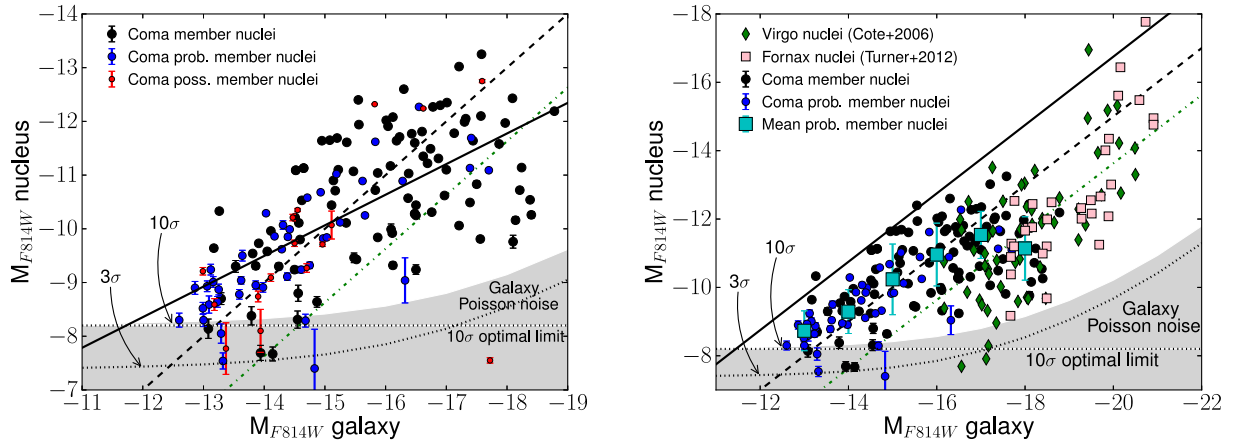
$$M_{\text{nuc}} = M_{\text{gal}} + (5.13 \pm 0.10) \quad (1)$$

$$M_{\text{nuc}} = M_{\text{gal}} + (5.54 \pm 0.10). \quad (2)$$

This is in agreement with earlier values (Graham & Guzmán 2003; Lotz et al. 2004; Grant et al. 2005, Table 1). Côté et al. speculate that the difference they find with earlier findings may be due to the use of different fitting functions or the higher sensitivity of their ACS observations, which allows them to detect fainter nuclei. Performing a fit with a free slope, we find

$$M_{\text{nuc}} = (0.57 \pm 0.05)(M_{\text{gal}} + 17.5) - (11.49 \pm 0.14), \quad (3)$$

where we assumed uniform error bars for all data points, since otherwise the fit was strongly biased by the few brightest NSCs in our sample. In fact, fitting the relation with a constant intrinsic Gaussian scatter gives an intrinsic dispersion of 0.9 mag, not very different from unit error bars. Our slope is marginally consistent (overlapping uncertainties) with the slope found by Graham and Guzman but less steep than what was found by Grant et al. The intercept is consistent with the value given by Graham & Guzmán (2003) but is about 0.5 mag off from the fitted value from Côté et al. Our slope agrees with those of Balcells et al. (2007,  $0.76 \pm 0.17$ ) and Scott & Graham (2013,  $0.60 \pm 0.10$ ).



**Figure 6.** Left: absolute  $F814W(AB)$  magnitudes of NSCs versus the absolute  $F814W(AB)$  magnitudes of the host galaxy for Coma members. The magnitudes were obtained by profile fitting. The host galaxy luminosity does not include the star cluster luminosity. The dashed line denotes our best-fitting line with unit slope. The green dash-dotted line represents a similar fit from Côté et al. The black solid line shows our best-fitting linear relation with non-unity slope. The grey areas approximate the regions where sky and galactic Poisson noise become dominant. Right: same plot, but now only containing members and probable members of the Coma cluster. Overplotted in green diamonds and pink squares are the uncorrected (see text)  $F850LP(AB)$  measurements from Côté et al. (2006) and Turner et al. (2012). We also overplot in big cyan squares the mean NSC magnitude and error on the mean by binning the likely members of Coma by host galaxy magnitude. The black solid line denotes the line where the NSC would contain 5 per cent of the light of the host galaxy.

**Table 1.** Luminosity scaling relations between NSCs and their host galaxies from this work and the literature.

Source	Relation	Band	Comments
Graham & Guzmán (2003)	$M_{\text{nuc}} = (0.87 \pm 0.26) (M_{\text{gal}} + 17.5) - (11.90 \pm 0.25)$	$F606W$	Outliers clipped
Grant et al. (2005)	$M_{\text{nuc}} = (0.74 \pm 0.06) (M_{\text{gal}} + 17.5) - (12.59)$	$I$ band	Zero-point estimated from plot
Côté et al. (2006)	$M_{\text{nuc}} = (1.0) (M_{\text{gal}} + 17.5) - (11.13 \pm 0.22)$	$F850LP$	Unit slope
	$M_{\text{nuc}} = (1.05 \pm 0.18) (M_{\text{gal}} + 17.5) - (11.05 \pm 2.19)$	$F850LP$	
Balcells et al. (2007)	$M_{\text{nuc}} = (0.76 \pm 0.17) (M_{\text{bulge}} + 25.0) - (15.5 \pm 0.45)$	$K$ band	For bulges
Scott & Graham (2013)	$M_{\text{nuc}} = (0.60 \pm 0.10) (M_{\text{gal}} + 20.4) - (16.57 \pm 0.175)$	$K$ band	
This work	$M_{\text{nuc}} = (0.57 \pm 0.05) (M_{\text{gal}} + 17.5) - (11.49 \pm 0.14)$	$F814W$	Uniform errors
	$M_{\text{nuc}} = (1.0) (M_{\text{gal}} + 17.5) - (12.36 \pm 0.10)$	$F814W$	

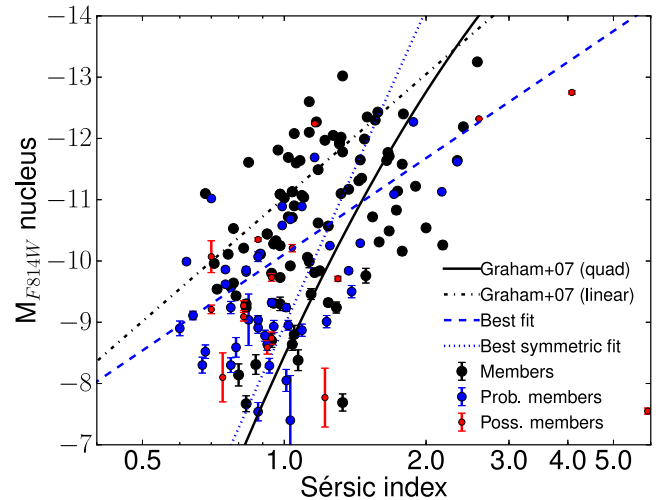
#### 4.4 Resolved star clusters

We find that a non-negligible number of star clusters ( $\sim 29$  per cent for the full sample, 25 per cent of all likely members) are resolved.

Although large (sizes up to 30 pc), we note that the sizes are not excessive: for example, the Milky Way GC NGC 2419 has a half-light radius of around 20 pc (Harris 1996; McLaughlin & van der Marel 2005), and also Mackey et al. (2006) find GCs around Andromeda with similar sizes, showing that these sizes are not uncommon for star clusters. We note that in principle it is possible that the objects with large sizes are not NSCs but nuclear discs. However, local NSCs are known to be discy (e.g. Seth et al. 2008). Since the difference between NSCs and nuclear discs is diffuse, we treat both resolved and unresolved sources in the same way when analysing scaling relations. In other words, we define an NSC to be the central excess light above a Sérsic profile.

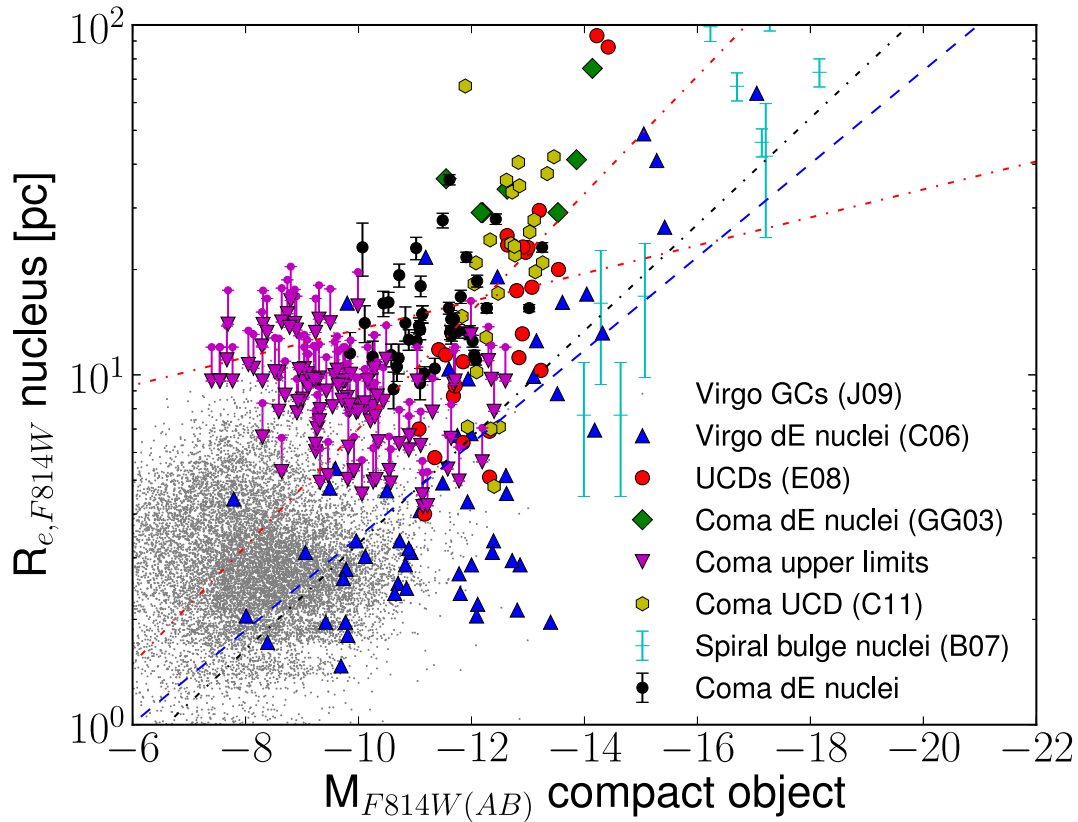
#### 4.5 Correlation with Sérsic index

In Fig. 7, we show the luminosity of the NSC as a function of the Sérsic index of the host galaxy. We note that we did not separate out different components (discs, bulges) so that the Sérsic index used here provides a measure of the overall concentration of the host galaxy. A comparison between the scatter on this relation with the scatter on the NSC/host galaxy luminosity relation (0.9 mag, see also Fig. 6) implies that the Sérsic index is not the main driver of



**Figure 7.** The luminosity of the NSC as a function of host galaxy Sérsic index. Overplotted are the curved and linear relations from Graham & Driver (2007) in black, which were converted to luminosity by assuming a constant mass-to-light ratio. Our own best-fitting relation  $M_{F814W} = (-5.2 \pm 0.7) \log_{10}(n/3) - (12.6 \pm 0.3)$  is given by the blue dashed line. The blue dotted line shows the result of a symmetrical regression fit, which minimizes the offsets perpendicular to the best-fitting line.





**Figure 8.** Effective radii of NSCs versus absolute NSC magnitudes in the  $F814W$  band (black circles) and upper limits (purple arrows, see text). Overplotted in blue triangles are the  $F850LP(AB)$  measurements from Côté et al. (2006), derived from Michie–King model fits to their data. Red circles are half-light radii of UCDs, from Evstigneeva et al. (2008), derived from fitting King models or combinations of King and Sérsic models. The green diamonds are NSCs in dEs in Coma fitted with a Gaussian component by Graham & Guzmán (2003). The grey small dots are the GC parameters from Jordán et al. (2009). The yellow hexagons are UCDs in the Coma cluster from Chiboucas et al. (2011), and the cyan pluses are the resolved nuclear components in spiral bulges from near-infrared photometry (Balcells et al. 2007). The dash-dotted (black) line is the relation of Bekki et al. (2004) and the double-dashed (blue) line shows the expected growth of NSCs by disposition of stars/gas at a radius determined by tidal forces. We fitted a linear relation to the sizes and magnitudes of resolved NSCs in Coma, both with and without taking into account upper limits on the unresolved sources (red dash-dotted lines).

NSC formation. Graham & Driver (2007) found a relation between the logarithm of the mass of an SMBH and the Sérsic index of the host galaxy with remarkably low scatter. They showed that this relation continues for low Sérsic index galaxies, when one replaces the mass of the SMBH with the NSC mass. They provide both a linear and curved CMO–Sérsic index relation. Graham (2012) argues that this relation should in fact be a broken relation. We convert the masses in the CMO-relations from Graham & Driver (2007) to luminosities and show these relations in Fig. 7, where we assumed a constant mass-to-light ratio in the  $I$  band of 2.0, based on Miles models (Vazdekis et al. 2010) for a 4 Gyr old stellar population with solar metallicity. A lower metallicity, which is likely for the fainter NSCs, will push these lines upwards, so that also the faintest NSC points will follow the curve on average. At a fixed Sérsic index, the points show significantly higher scatter than the 0.9 mag in the nucleus–host galaxy luminosity scaling relation. If indeed NSCs form the counterpart of SMBHs at low galaxy masses – and this is debated (e.g. Graham 2012) – it remains to be explained why the scatter in NSC magnitudes is so much higher than for the  $M_{\text{BH}}-n$  relation. For completeness, we note that our best-fitting relation between magnitude and Sérsic index is  $M_{F814W} = (-5.2 \pm 0.7)\log_{10}(n/3) - (12.6 \pm 0.3)$ . A symmetric fit that minimizes the distances perpendicular to the best-fitting line yields  $M_{F814W} = (-17.2 \pm 4.7)\log_{10}(n/3) - (17.1 \pm 3.9)$ , which has a

slope that is similar to the log-quadratic fit of Graham at low Sérsic indices. A log-quadratic fit is not warranted by our data.

Emsellem & van de Ven (2008) argue that the tidal field in galaxy centres becomes disruptive for steep inner profiles, and that one therefore should not expect to see many NSCs in galaxies with Sérsic index  $n \gtrsim 3.5$ . Since we selected dEs, which preferentially have low Sérsic indices, we are not able to check this prediction. However, trends with Sérsic index are expected, and we will discuss this more in Section 5.2.

## 5 DISCUSSION

### 5.1 Sizes of nuclear clusters

In Fig. 8, we show sizes of NSCs as a function of the NSC’s absolute magnitude, together with similar measurements of NSCs (from the ACS VCS and from near-infrared photometry of nearby spiral bulges), UCDs and GCs. The measured sizes of Coma NSCs are similar to sizes of the other compact systems. Unresolved NSCs are shown as upper limits. Since our ability to resolve NSCs depends on S/N, we use upper limits from the BAGATELLE fits of resolved clusters to the data, where we define an upper limit as that size for which we know with 67 per cent certainty that the cluster is smaller than that. We have limited our analysis to dEs and have not analysed the

bright nuclear excess light typically seen in  $L_*$ -ellipticals, as was done for the ACS VCS. The Coma NSC points cover a different locus of the luminosity–size diagram than the NSCs from the ACS VCS, in the sense that at a fixed luminosity NSCs in Coma are somewhat larger. We find that  $\sim 70$  per cent of star clusters with  $F814W$  absolute magnitude between  $-14$  and  $-11$  have half-light radii  $>9$  pc, whereas for the Virgo ACS this percentage is closer to 25 per cent.

A possible reason for this may be the use of different functional forms for fitting the NSC (we use a Gaussian profile; Côté et al. use a King profile) as well as the use of different PSFs. We have tested the influence of a different functional form by fitting a source with an empirical King profile and find that the effective radius is consistent with the one inferred from a Gaussian, unless we put strong constraints on the concentration. More likely is that, despite all our tests to confirm the sizes of the NSCs, we may be overestimating the quality of our data or our knowledge about the PSF, or, alternatively, the structural parameters fits may be biased to resolved sources.

Both the upper-right part of the diagram and the lower-left part are uninhabited: there are no luminous NSCs with small sizes, or faint NSCs with large sizes. Although the latter could be caused by selection effects, the first relation probably reflects the formation of bright NSCs.

If the sizes of the Coma NSCs are indeed larger than the sizes of nuclear clusters of Virgo dwarfs, how can this be explained? Given the compactness of NSCs in general, it is hard to believe that tidal interactions with other galaxies in the cluster have modified the structure of the NSC. However, the NSCs in Coma and Virgo may be at different evolutionary stages; it is possible that some of the secular processes mentioned in the introduction (absorption of kinetic energy from the host galaxy, puffing up by a central BH) are responsible for the larger sizes.

The GC merging scenario for NSCs gives predictions for the sizes of NSCs as a function of mass (e.g. Bekki et al. 2004). We can infer the slope of the mass–size relation based on just our Coma data. We fit a linear relation between magnitude and size to the resolved clusters in our sample. Under the assumption of a constant mass-to-light ratio, this relation is

$$\text{Size} \propto \text{Mass}^{0.1 \pm 0.05}. \quad (4)$$

This relation does not take into account that for most of the sources in our sample we only have upper limits on the size. We know the luminosity of these sources. However, the size can be anything between roughly 1 and 8 pc. Adding the unresolved sources to the likelihood, we find that the sizes of the Coma NSCs follow a slightly steeper relation:

$$\text{Size} \propto \text{Mass}^{0.42 \pm 0.04}. \quad (5)$$

In Fig. 8, we also show the size–luminosity relation predicted by Bekki et al. (2004) for NSCs formed from merging GCs. The slope of this relation is consistent with the one of the Coma NSCs including upperlimits.

For the dissipational formation of NSCs, there exist no predictions for a mass–radius relation. However, Seth et al. (2006) discuss the formation of a stellar ring around the NSC in IC5052 from the tidal disruption of a molecular cloud by the already existing NSC. If this mechanism is universal and NSCs grow inside-out with new material deposited at a radius determined by tidal forces, we expect a mass–size relation with the same slope as the mass–size relation predicted by Antonini (2013) for GC accretion ( $\text{Size} \propto \text{Mass}^{1/3}$ ).

Several authors have compared sizes and luminosities of NSCs with those of UCDs. We note that in the same luminosity range as the UCDs of Chiboucas et al. (2011), the NSCs and UCDs in Coma dEs have similar sizes. If UCDs are the high-mass end of the GC luminosity function, this similarity in size is not unexpected, since the high-mass GCs are the ones that most easily reach the centre. Similarly, if UCDs are stripped dEs, the sizes of UCDs and dE NSCs should be similar. Given the overlap in luminosity between NSCs and GCs (Fig. 5) a part of the GC cluster population brighter than  $F814W = -10$  mag likely consists partly of stripped NSCs.

## 5.2 Origin of the $L_{\text{nuc}}-L_{\text{gal}}$ relation

There exist predictions for the scaling relation between the mass of the host galaxy and the mass of the NSC under the assumption that it formed entirely from GC accretion. Antonini (2013) provides predictions and derivations for both. For GC accretion, the mass of the star cluster is

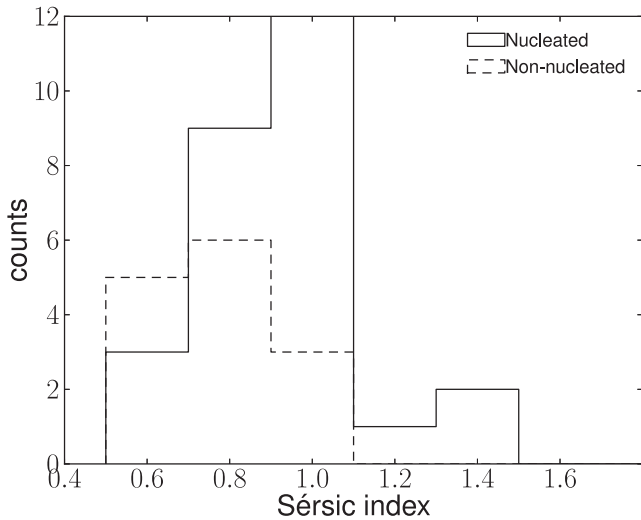
$$M_{\text{NSC}} = 3 \times 10^7 \left( \frac{f}{10^{-6}} \right) \left( \frac{\ln \Lambda}{3} \right) \left( \frac{m_*}{M_\odot} \right) \times \left( \frac{\langle m_{\text{cl}} \rangle}{10^5 M_\odot} \right)^{3/2} \left( \frac{t}{10^{10}} \right)^{1/2} \left( \frac{\sigma}{50 \text{ km s}^{-1}} \right)^{3/2}$$

with  $f$  the initial number fraction of GCs and  $\ln \Lambda$  the Coulomb logarithm, which is defined as  $\Lambda = \frac{b_{\text{max}} v_*^2}{G(M_{\text{GC}} + m_*)}$ , with  $b_{\text{max}}$  the largest possible impact parameter and  $m_*$  and  $v_*$  the typical masses and velocities of the stars.

Since for low-mass galaxies the Faber–Jackson relation scales as  $L \propto \sigma^2$  (Davies et al. 1983; Held et al. 1992; Matković & Guzmán 2005), this yields a prediction between the magnitude of the star cluster and the host galaxy with slope 0.75 (with maybe a small additional dependence on the GC fraction with host galaxy magnitude). For a more complex model of GC accretion, Gnedin, Ostriker & Tremaine (2014) derive that the mass of the NSC and the mass of the galactic spheroid scale as  $M_{\text{NSC}}/M_{\text{star}} \approx 0.0025 M_{\text{star},11}^{-0.5}$ , i.e. with a slope of 0.5. Here,  $M_{\text{star},11}$  denotes the stellar mass of the spheroid divided by  $10^{11} M_\odot$ .

For *in situ* formation of the cluster, the power-law slope between  $M_{\text{NSC}}$  and  $\sigma$  is predicted to be 4.0 for a constant momentum feedback model (McLaughlin et al. 2006; Antonini 2013), leading to a significantly steeper slope between NSC magnitude and host galaxy magnitude of 2.0. This slope seems to be consistent with the Virgo and Fornax data in Fig. 6. However, it is significantly steeper than our fitted slope.

Although a correlation between  $\sigma$  and NSC mass with slope 4.0 was reported by Ferrarese et al. (2006), following up on this work, Scott & Graham (2013) found a much shallower slope with  $M_{\text{NSC}} \propto \sigma^{2.11}$ . Our slope seems to be consistent with that. Although the coincidence of the slopes for the GC accretion scenario and the slope found above are striking, it remains to be seen if dissipational formation of NSCs is ruled out by this. Since the momentum feedback in the McLaughlin model may not be important at later times, when the winds from supernovae and young stars have faded, the predicted slope of 4 may be overestimated compared to reality. As an example, Seth et al. (2006) came up with a model in which gas is accreted episodically in the centre of the host galaxy. Although it depends critically on the gas physics, it is not unthinkable that if a fixed fraction of the gas in the galaxy ends up in the nucleus, the NSC mass grows proportionally to the mass of the host galaxy. We conclude that, although the slope of the host galaxy–NSC luminosity relation is in good agreement with the GC inspiral scenario,



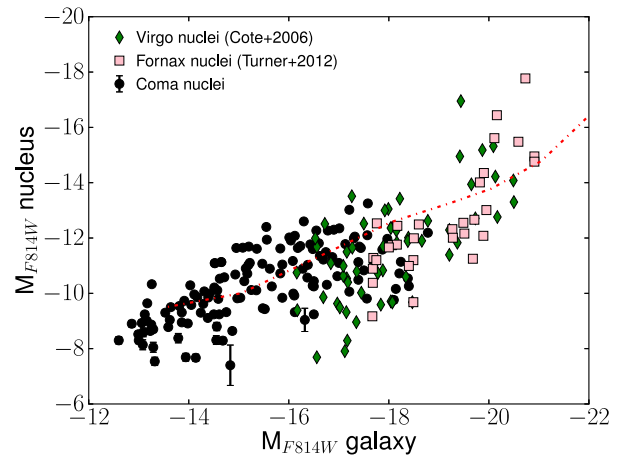
**Figure 9.** Histograms of the distribution of Sérsic indices for nucleated (solid line) and non-nucleated (dashed line) galaxies fainter than  $M_{F814W} = -14$ . The mean Sérsic index of non-nucleated galaxies is slightly, but not significantly, lower than for nucleated galaxies ( $0.76 \pm 0.17$  versus  $0.92 \pm 0.18$ ), which implies that the missing central star cluster in non-nucleated galaxies is not mistaken for a galaxy with a more cusped profile.

the lack of detailed model predictions for the formation and evolution of NCS through dissipational collapse make the exclusion of the latter scenario premature. It may be possible that non-nucleated galaxies do contain a NSC which is too faint to be detected in our data. As a sanity check, we have compared for the lowest luminosity galaxies ( $M_{F814W} < -14$ ) the Sérsic indices of non-nucleated galaxies with those of nucleated galaxies. The Sérsic indices of non-nucleated galaxies are on average lower than those of nucleated galaxies, suggesting that at least the nucleus is not mistaken for a central cusp (Fig. 9). We also checked by how much the scaling relations would change if we would replace the non-detections by upper limits. The results suggest that the main change happens in the intercept of the relation and the intrinsic scatter, and not so much in the slope (which changes by less than 0.1). If we allow for a variable nucleation fraction among the non-detections, 10 per cent of the non-detections would be consistent with our fitted relation. We note that although this is comforting, it is by no means solid evidence that these galaxies are not nucleated.

In general, we find that NSCs are almost never brighter than 5 per cent of the host galaxy. This is slightly different for the bright galaxies in the Virgo and Fornax samples; however, we will argue that structurally (in this paper) and in terms of stellar populations (in a follow-up paper) the NSCs in those galaxies have probably formed or evolved through a different channel. Galaxies with  $F814W$  magnitude  $\sim -18$  appear to not reach this maximum NSC formation efficiency. It is well known that the specific frequency of GCs is curved in this magnitude range (e.g. Peng et al. 2008) – as the formation of GCs is happening mainly in the very early Universe before feedback shuts off the star formation in the host galaxy, the GC specific frequency can be high for low-mass galaxies. This may be a possible explanation for the curvature in the NSC–host galaxy magnitude relation.

### 5.2.1 Dependence on Sérsic index

Both formation scenarios predict that the mass of the star cluster is proportional to the mass of the host galaxy. For dissipational



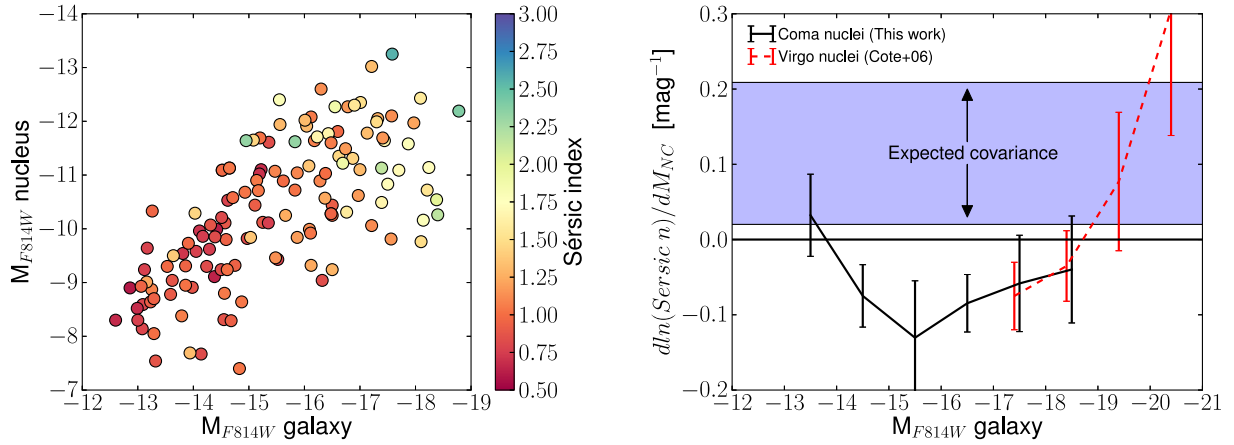
**Figure 10.** Magnitude of NSCs versus host galaxy magnitude, with the average luminosity in gGCs (red dash-dotted line), based on measurements in the Virgo cluster.

collapse, a larger gas reservoir has been available for formation of the nuclear cluster. For the GC accretion scenario, the number of GCs is proportional (although not linearly) to the mass of the host galaxy. We now try to distinguish between these two scenarios by looking at the scaling relations of NSCs and host galaxies.

Fig. 10 shows the luminosity–luminosity relation with superimposed the average luminosity of GC systems of dEs in Virgo, based on the data of Peng et al. and Lotz et al. There is a strange coincidence, which has been noted before (e.g. Côté et al. 2006), that the luminosity of the NSCs is similar to the combined luminosity of all GCs in the galaxy. If NSCs form purely from accretion of GCs, this means that GC inspiralling in low-mass galaxies is less efficient. The NSCs in high-mass galaxies stand out as well, as they do not follow the GC luminosity line.

In Fig. 11, we show how the NSC luminosity versus host galaxy luminosity depends on the Sérsic index. In this plot we show, for a given magnitude bin, the slope of a linear relation between NSC luminosities and Sérsic indices (assuming unit error bars while fitting). For the low-mass end ( $M_{F814W} \sim -13$ ), this relation may be influenced by the low S/N of the data, however, for intermediate-mass galaxies, we find that in a given magnitude bin, galaxies with high Sérsic indices have on average higher NSC luminosities.. We also determine this dependence for galaxies in the Virgo ACS survey, and find that the relations overlap, in the region where the samples overlap.

The Sérsic index dependence at the low-mass end was found to be different from the slope found for intermediate-luminosity dEs. However, the Sérsic index and NSC luminosity of a galaxy are expected to be covariant: if we fit to a galaxy a profile with a Sérsic index that is slightly lower than the actual Sérsic index (*ceteris paribus*), thereby lowering  $\mu_0$ , we may still be able to obtain a good fit if we compensate the missing central light by increasing the NSC luminosity. As our fitting code explores different parameter configurations and calculates a posterior probability distribution for each of them, we can infer the magnitude of this effect by diagonalizing the covariance matrix for Sérsic index and NSC magnitude. Since it is difficult to calculate the uncertainty on the covariance from the data of a single galaxy (as this would involve the calculation of an additional covariance matrix for the elements of the covariance matrix), we take the standard deviation of the values found for several galaxies. From the Markov chain Monte Carlo (MCMC) output of our fitting code, we thus determine that the average covariance between



**Figure 11.** Left: magnitude of NSCs versus host galaxy magnitude, colour-coded by host galaxy Sérsic index. Right: the dependence of NSC luminosity on Sérsic index, for different host galaxy magnitudes, for both the Coma likely members sample (solid black line) and the ACS Virgo sample (dashed red line). Negative values mean that a higher Sérsic index implies a more luminous NSC, contrary to the covariance between NSC luminosity and Sérsic index determined from the MCMC output (purple region, see text).

Sérsic  $n$  and  $M_{\text{nuc}}$  translates into  $d \ln n / d \text{mag} = 0.114 \pm 0.094$ . This is close to the value that we find for the magnitude  $-13$  bin, suggesting that that particular value may be due to low S/N. More importantly however, the sign of the slope inferred from the covariance is opposite the sign of the result for intermediate-luminosity dEs, as already qualitatively argued above, suggesting that this result is solid.

We have thus found that galaxies with a relatively high Sérsic index have on average a more luminous NSC. The explanation may be two-fold: galaxies with high Sérsic index may produce more gas in the centre from stellar outflows, and may also be better at retaining gas since the potential is deeper, which may also affect the ambient pressure. On the other hand, GC accretion may be more efficient in galaxies with higher Sérsic index. A test of the first scenario would require hydrodynamic simulations, which go beyond the scope of this paper. We do note that this trend is opposite what is expected from tidal forces in galaxy centres (Emsellem & van de Ven 2008). In order to see if we can reproduce the Sérsic index dependence with the GC accretion scenario, we make a simple model for the production of NSCs.

### 5.2.2 A toy model for GC accretion

This model<sup>3</sup> for the accretion of GCs is based on the dynamical friction formula of Chandrasekhar (1943) for a Gaussian velocity distribution (note that we make strong assumptions on the geometry here, namely sphericity and the size scales which are taken into account by the Coulomb logarithm):

$$f = -4\pi G^2 M_{\text{GC}} \rho(r) \ln \Lambda \left[ \text{erf}(X) - \frac{2X}{\sqrt{\pi}} e^{-X^2} \right] \frac{v}{v^3}, \quad (6)$$

where  $v$  is the velocity of the inspiralling GC and  $X = v/(\sqrt{2}\sigma)$ , with  $\sigma$  the local velocity dispersion. Since we have no handle on the actual value of  $\ln \Lambda$ , the Coulomb logarithm, we make the somewhat ad hoc assumption that  $\ln \Lambda = 3$  for all systems. Spinnato et al. (2003) find that  $\ln \Lambda = 2.9$  for dense star clusters, and 6.6 for

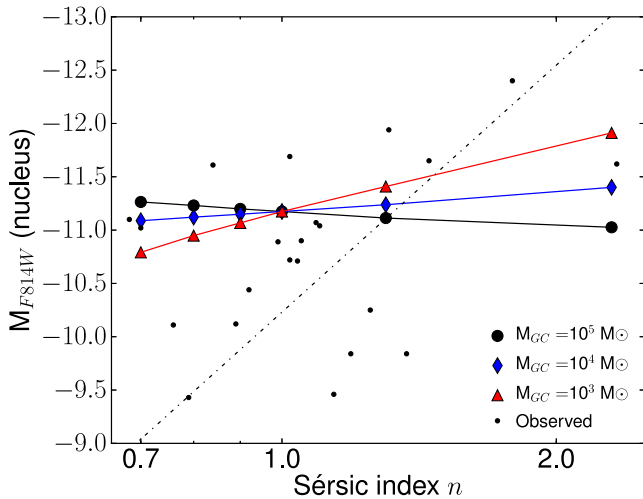
massive BHs. We note however that even assuming a value for the Coulomb logarithm as high as  $\ln \Lambda = 10$  does not change the conclusions. We assume that all matter follows visible matter, and that the projected mass–density is therefore described by a Sérsic profile. In the model, GCs are orbiting the galaxy centre with the circular velocity at each galactocentric radius, which we calculate from the total enclosed mass by integrating the deprojected Sérsic profile. The velocity dispersion at each radius is calculated from the spherical Jeans equation, assuming isotropy throughout the galaxy. We thus calculate how much time it takes for a GC with mass  $M_{\text{GC}}$ , starting out at radius  $r$ , to reach the centre of the galaxy.

This allows us to calculate the maximum radius for which a GC with mass  $M_{\text{GC}}$  can reach the galaxy centre in less than a Hubble time. Given this radius, we assume that a fixed fraction of the enclosed mass in the galaxy formed in GCs and produces the NSC. Leaving all parameters (mass, effective radius) for the model galaxy fixed, except the Sérsic index, we then determine how a change in Sérsic index changes the NSC mass.

Our model suggests that if an NSC forms from low-mass GCs ( $10^3 M_{\odot}$ ), then only GCs that are already close to the galaxy centre manage to reach the centre in a Hubble time, since the dynamical friction times for low-mass GCs in the outer part of the galaxy are too long. On the other hand, high-mass GCs ( $M = 10^5$ – $10^6 M_{\odot}$ ) may form farther out than low-mass GCs and still reach the centre, since the dynamical friction time is proportional  $M^{-2}$ . Massive clusters, which start their journey to the galaxy centre at a distance of a few effective radii, in a galaxy for which we have increased the Sérsic index slightly, reach the centre somewhat quicker indeed. However, this shorter inspiralling time and therefore larger initial radius does not lead to a larger NSC mass, since the increased Sérsic index redistributes the mass in the galaxy in such a way that there is more mass in the outer parts, and therefore the increase in  $n$  does not lead to an increase in the enclosed mass within the maximum inspiral radius of the GC. This is illustrated in Fig. 12 for a  $10^5 M_{\odot}$  GC, where we assumed a model dwarf galaxy with  $M_{\text{F814W}} = -15$ , an effective radius of 700 pc and  $M/L = 2.3$  in the  $I$  band and calculate for five different Sérsic indices the expected NSC mass, which we normalized at  $n = 1$ . Not only is the dependence on Sérsic index much weaker than expected, also the sign of the relation is wrong for high-mass GCs. For low-mass GCs ( $10^3 M_{\odot}$ ), we find that the

<sup>3</sup> We note that a similar calculation can be found in section 8.2.2 of Merritt (2013).





**Figure 12.** The final magnitude of the NSC after inspiralling of different mass GCs for a Hubble time for different Sérsic indices of the host galaxy. The host galaxy is assumed to be an  $F814W = -15$  spherical dE with  $M/L = 2.3$ . The low-mass GCs have to form close to the centre to build up the NSC in a Hubble time and therefore show a strong dependence on Sérsic index. The high-mass GCs can form farther out for higher Sérsic indices, but in these galaxies the mass is also distributed in such a way that Sérsic index dependence actually changes sign. The small dots show the observed probable member data in the  $-16 < F814W < -15$  bin, together with the fit assuming uniform errors on  $\log n$ .

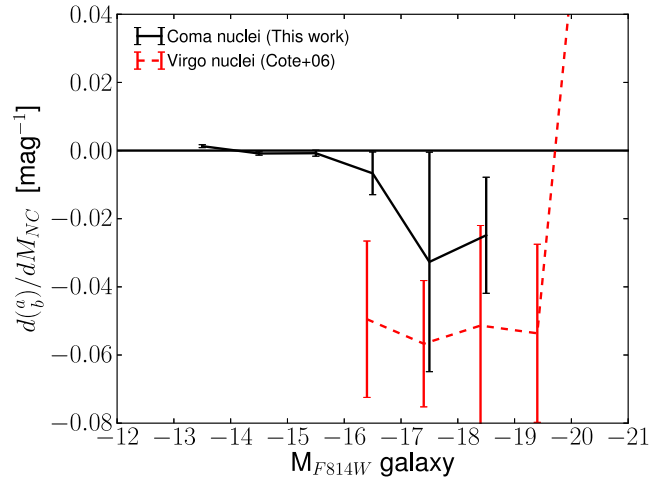
trend with Sérsic index has the right sign, but is not strong enough to explain the observed trend.

Our model is based on strong assumptions (Sérsic index constant over time, fixed fraction of stars form in clusters of similar mass) but at face value suggests that we cannot reproduce the Sérsic index–luminosity trend unless we build up our star clusters from low-mass building blocks close to the galaxy centre. A more natural explanation may however be a small amount of residual dissipational star formation.

### 5.2.3 Dependence on host galaxy flattening

We note that other authors (van den Bergh 1986; Ryden & Terndrup 1994) have already found that non-nucleated galaxies are generally more elliptical than nucleated galaxies. Our data confirm, that, for galaxies with similar luminosities, non-nucleated galaxies tend to have a higher axis-ratio than nucleated galaxies.

However, not only the nucleation fraction, but also the host galaxy–NSC luminosity appears to depend on the ellipticity of the host galaxy. In Fig. 13, we show a relation for the slope of the axis ratio–NSC luminosity as a function of host galaxy magnitude. For luminous dEs, the rounder galaxies tend to have brighter NSCs. A simple explanation for this may be that dEs that are rounder have been in the cluster for a longer time and hence have lost some of their mass, that is, instead of the NSC becoming brighter over time, the host galaxy has become fainter. It is unclear if the GC accretion scenario predicts this scaling with axis ratio, because, although dynamical friction may be less efficient outside the plane of the dEs, most globulars actually seem to lie in the plane of the galaxy (Beasley et al. 2009; Wang et al. 2013). For the dissipational model, it is possible that galaxies that are rounder have less rotational support, making it easier for the gas to reach the centre.



**Figure 13.** The slope of the NSC luminosity–host galaxy axis ratio relation as a function of host galaxy magnitude.

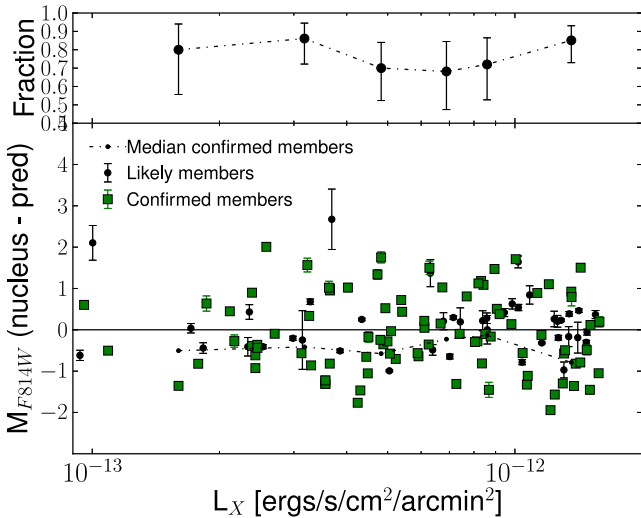
Interestingly, Seth et al. (2008) find that late-type spiral galaxies have NSCs that are on average one order of magnitude less massive than those found in elliptical galaxies of the same mass. An obvious explanation is that late-type spirals are the progenitors of early-type dEs, which may have lost part of their outer parts due to tidal stripping, or alternatively, have grown their NSC disproportionately compared to their host galaxy after entering the cluster. It may thus be the case that this relation tells us more about the evolution of the host galaxy in the cluster environment than about the formation of the NSC.

### 5.3 Influence of environment

If nuclei are formed or still evolve after a galaxy falls into the cluster, one might expect that properties of the nuclei depend on environment. As the pressure of the intracluster medium (ICM) increases towards the centre of the cluster, dwarf galaxies near the cluster centre may have had additional bursts of star formation, because the ICM confines the gas (Babul & Rees 1992).

In Fig. 14, we show the difference in NSC magnitude and host galaxy magnitude versus the X-ray flux at the centre of the host galaxy. For this, we used the *XMM* 0.5–20 keV map of Finoguenov, Briel & Henry (2003), but note that using clustercentric distance does not change the results. If the hot cluster gas plays any role in enhancing the evolution of NSCs, the difference in magnitude between NSC and host galaxy should be higher at high X-ray flux. We find no indication that this is the case. We bin together data points in bin width of  $2 \times 10^{-13} \text{ ergs s}^{-1} \text{ cm}^{-2} \text{ arcmin}^{-2}$  and calculate for each bin the median difference between the NSC and the host galaxy. If anything, the magnitude difference is slightly higher in high-density regions. An analysis for different host galaxy magnitude bins, similar to what we did in Section 5.2 for Sérsic index and ellipticity, does not reveal any environmental trend either.

The virial radius of the Coma cluster is approximately 2.9 Mpc (Łokas & Mamon 2003). Taking  $\langle v_{\text{rms}} \rangle \approx 1000 \text{ km s}^{-1}$  as a typical velocity for a dwarf galaxy in the cluster (Struble & Rood 1999), a crossing time is approximately 3 Gyr. If we assume that dwarf galaxies have ages in the range 4–7 Gyr (Koleva et al. 2009; Smith et al. 2009) and that the cluster has been responsible for shutting off star formation as soon as a dwarf passed the virial radius, a typical dE has already passed the cluster once or twice, and (depending



**Figure 14.** Upper panel: dependence of nucleation fraction on environment, as parametrized by the 0.5–20 keV X-ray flux. The error bars denote the 95 per cent Jeffreys interval. Lower panel: difference between NSC magnitude and average NSC magnitude for a fixed host galaxy magnitude bin as calculated in Fig. 5, shown as a function of 0.5–20 keV X-ray flux.

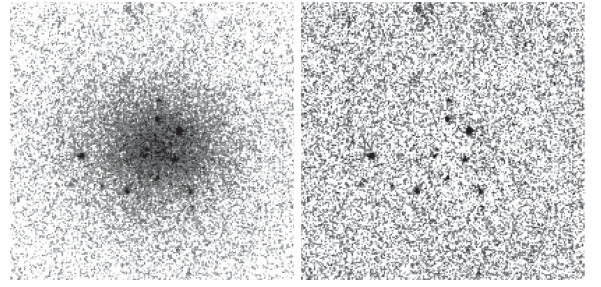
on the orbit) may therefore have seen different environments of the cluster. It is therefore possible that the influence that environment may have had on the formation or evolution of the NSC is not reflected by the current environment of the dwarf galaxy.

Lisker et al. (2007) analysed the clustering and flattening of nucleated dEs in the Virgo cluster, and found that these galaxies, compared to non-nucleated and disk dEs, formed a more relaxed population. This is consistent with the lack of trend with environment that we see in Coma. The nucleation fraction does suggest a change with density – although the statistics are poor – but not in a monotonic way. It is possible that nucleation is a secular process and that instead the formation/destruction of non-nucleated galaxies is dependent on environment. As an example we mention the possibility that some of the non-nucleated galaxies form as gas-poor tidal galaxies, and therefore should show a dependence on environment (e.g. Okazaki & Taniguchi 2000; Chattopadhyay et al. 2015). Since the non-nucleated galaxies have on average lower Sérsic indices, they may be more easily destroyed in the cluster centre by tidal forces than the more centrally concentrated nucleated galaxies.

#### 5.4 Non-nucleated galaxies

In the left-hand panel of Fig. 4, we showed how the detection fraction of NSCs depends on magnitude. The absence of NSCs in low-mass galaxies has been known since a long time (van den Bergh 1986). Turner et al. (2012) suggest that galaxies may not be able to form nuclei when the NSC magnitude would fall below the turnover magnitude of the GC luminosity function.

Several galaxies do not show an NSC, although in a few cases, there are GCs surrounding the galaxy. As an example, we show in Fig. 15 the galaxy SDSSJ125636.63–271503.6, which, since it is in class 2, belongs to Coma with a 90 per cent probability. The galaxy is not nucleated, but shows several clusters within one effective radius, though all of them several pixels away from the photometric centre. The spatial coincidence of these clusters and the galaxy implies that they are physically associated.



**Figure 15.** The non-nucleated galaxy SDSSJ125636.63–271503.6 and the residual from the BAGETTE fit. Several star clusters are visible, but none of them is at the photometric centre of the galaxy.

The galaxy, found in one of the more remote tiles (visit 63), is rather faint ( $F814W = -15.9$ ), and the colour, which compared to other galaxies in the sample is typical for a galaxy of this magnitude, implies a single stellar population (SSP) age (assuming  $[Z/H] = -0.3$ , half the solar metallicity) of more than 3 Gyr. The core of this galaxy appears to be slightly redder than the outer parts, with no indication of the bluer central colour often found for NSCs. The brightest six clusters all have  $F814W$  magnitudes between 24.4 and 25.9. If we exclude the brightest and the faintest cluster, the four remaining clusters have a combined absolute magnitude of  $-11.4$ , consistent with the galaxy luminosity–NSC luminosity relation.

This galaxy poses a problem for both formation scenarios: if GC inspiralling is the dominant mechanism for NSC formation, why is there no NSC in the centre of this galaxy, since the most massive clusters should have had plenty of time to reach the centre? And if a galaxy is able to form GCs in its outer parts, why is it not able to form an NSC in its centre, where the high pressure is probably favourable for the formation of a huge star cluster? This galaxy shows that, even though the ability to form GC-mass objects may be a necessity to form an NSC, it is apparently not always sufficient.

## 6 CONCLUSIONS

We have analysed data of a sample of 200 dE galaxies in the Coma cluster, to study the scaling relations of their NSCs, and in doing so, double the sample of *HST*-observed NSCs in dEs. Our conclusions are the following.

- (i) NSCs are present in almost all low-mass galaxies. The detection fraction is consistent with what has been found for Virgo by Côté et al. and previously for Coma by Graham & Guzman.
- (ii) The luminosity function of NSCs peaks 1.5 mag fainter than the luminosity function of the Virgo sample of Côté et al. A natural explanation is the difference in host galaxy magnitudes between the samples. This also means that the luminosity function is closer to the peak of the GC luminosity function.
- (iii) The magnitudes of NSCs follow a curved relation w.r.t. host galaxy magnitude. A possible explanation for the shape at the bright end is that the excess light is due to wet or moist mergers. At the faint end star cluster formation may have been less efficient due to feedback or long inspiralling time-scales. A linear fit between the NSC and host galaxy magnitude gives  $M_{\text{nuc}} = (0.57 \pm 0.05)(M_{\text{gal}} + 17.5) - (11.49 \pm 0.14)$ .
- (iv) At fixed magnitude, galaxies with higher Sérsic indices tend to have brighter NSCs. A toy model for the formation of NSCs by the inspiralling of GCs does not predict the right trend or order of magnitude, although we cannot exclude that more sophisticated

modelling will do so. A plausible explanation is, that additional dissipational star formation in the centre of galaxies with high Sérsic indices is more efficient since the gas is more easily retained.

(v) Rounder galaxies have slightly more luminous NSCs. The most obvious explanation is that these galaxies have been in the cluster for longer time and have lost part of their mass. On the other hand, it is possible that dynamical friction is less efficient in flattened galaxies or that angular momentum prevents gas to reach the centre.

From the scaling relations of NSCs a picture appears in which their formation is consistent with the GC accretion scenario. Despite this, the GC accretion model still fails to explain why the nucleation fraction goes down at the faint end where dynamical friction should be more efficient, and similarly does not explain the trend with Sérsic index.

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## APPENDIX A: TABLES AND PROFILES

## APPENDIX B: FITTING CODE

We make use of our own custom fitting code, called `BAGATELLE`, to measure the structural parameters of nuclei and host galaxies. Our code is a two-dimensional fitting code, similar to `GALFIT` (Peng et al. 2002), but based on a Bayesian framework for exploring the posterior distribution of the fitted parameters. It therefore allows

comparison between simple and complex models (e.g. with and without nuclear cluster) and accurate estimates of covariances between parameters and more accurate error bars. It is thus similar to the `GALPHAT` code of Yoon, Weinberg & Katz (2011).

## B1 Model generation

The basis for each model is a one-dimensional surface brightness profile, for which we adopted variations of the Sérsic profile (see equation 1). We are forced to choose priors for each free parameter in the fit. These are not hard-coded, but we list them here. The priors on the host galaxy magnitude and effective radius are based on the output of `SEXTRACTOR` (we choose a 4.0 mag interval around the `SEXTRACTOR` magnitude and a factor of 15 interval around the `SEXTRACTOR` radius). The Sérsic index is allowed to vary between 0.3 and 10 for all galaxies. The nuclear cluster magnitude is allowed to vary between the lower limit of the host galaxy magnitude and the detection limit of our data. For Gaussian sources, we allowed the full width at half-maximum (FWHM) to vary between 0.4 and 15 pixels. We note that for good-enough data, the output parameters are basically insensitive to the choice of prior.

Each two-dimensional coordinate ( $x, y$ ) corresponds to a surface brightness, where we assume that all isophotes are elliptical – i.e. we take into account the ellipticity and position angle of the model, but do not make use of generalized ellipses. The model centre falls exactly on a pixel centre (subpixel shifts of the model are treated during the convolution), and each pixel is subsampled by a factor of  $20 \times 20$ , which is increased in the centre of the galaxy to a factor over 100, depending on the steepness of the profile.

## B2 PSF convolution

The model is then convolved by the PSF. For this, we first convolve the PSF with a two-dimensional modulated sinc function (to account for subpixel shifts), and then convolve the model with it. Convolution is done in the Fourier domain. We use the fast Fourier transform

**Table A1.** Structural parameters of galaxies used in the paper. The complete table is available in the online version of this paper.

Name	RA	Dec.	$P_m$	$F814W$ (gal)	Sérsic $n$	$F814W$ (nuc)	FWHM
SDSSJ130018.54–280549.7	195.0773	28.0972	0	$16.8 \pm 0.01$	$1.5 \pm 0.01$	$24.3 \pm 0.05$	–
LEDA126789	194.8829	27.8613	0	$16.9 \pm 0.01$	$1.6 \pm 0.01$	$22.6 \pm 0.03$	$2.3 \pm 0.1$
SDSSJ130041.19–280242.4	195.1717	28.0451	0	$16.8 \pm 0.01$	$1.7 \pm 0.01$	$23.9 \pm 0.04$	–
SDSSJ130034.42–275604.9	195.1435	27.9347	0	$16.6 \pm 0.01$	$2.2 \pm 0.01$	$24.7 \pm 0.06$	–
RB068	194.9978	27.9406	0	$16.2 \pm 0.01$	$2.4 \pm 0.01$	$22.8 \pm 0.01$	–
LEDA126815	194.6898	27.7539	0	$16.62 \pm 0.01$	$2.0 \pm 0.01$	$24.46 \pm 0.04$	–
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
⋮	⋮	⋮	⋮	⋮	⋮	⋮	⋮
COMAi125932.883p275800.05	194.8870	27.9666	2	$21.84 \pm 0.36$	$1.2 \pm 0.3$	$26.0 \pm 0.1$	–
COMAi125929.995p275348.12	194.8748	27.8967	2	$21.14 \pm 0.17$	$1.0 \pm 0.1$	$26.0 \pm 0.1$	–
COMAi125952.543p275824.21	194.9688	27.9734	2	$21.94 \pm 0.22$	$0.9 \pm 0.2$	$26.1 \pm 0.1$	–
COMAi13021.712p275650.16	195.0904	27.9472	2	$21.41 \pm 0.16$	$0.9 \pm 0.1$	$26.2 \pm 0.1$	–
COMAi125944.017p275615.29	194.9333	27.9375	0	$21.71 \pm 0.06$	$0.9 \pm 0.1$	$26.3 \pm 0.1$	–
COMAi125924.938p275320.35	194.8538	27.8889	2	$21.71 \pm 0.13$	$1.0 \pm 0.1$	$27.0 \pm 0.2$	–

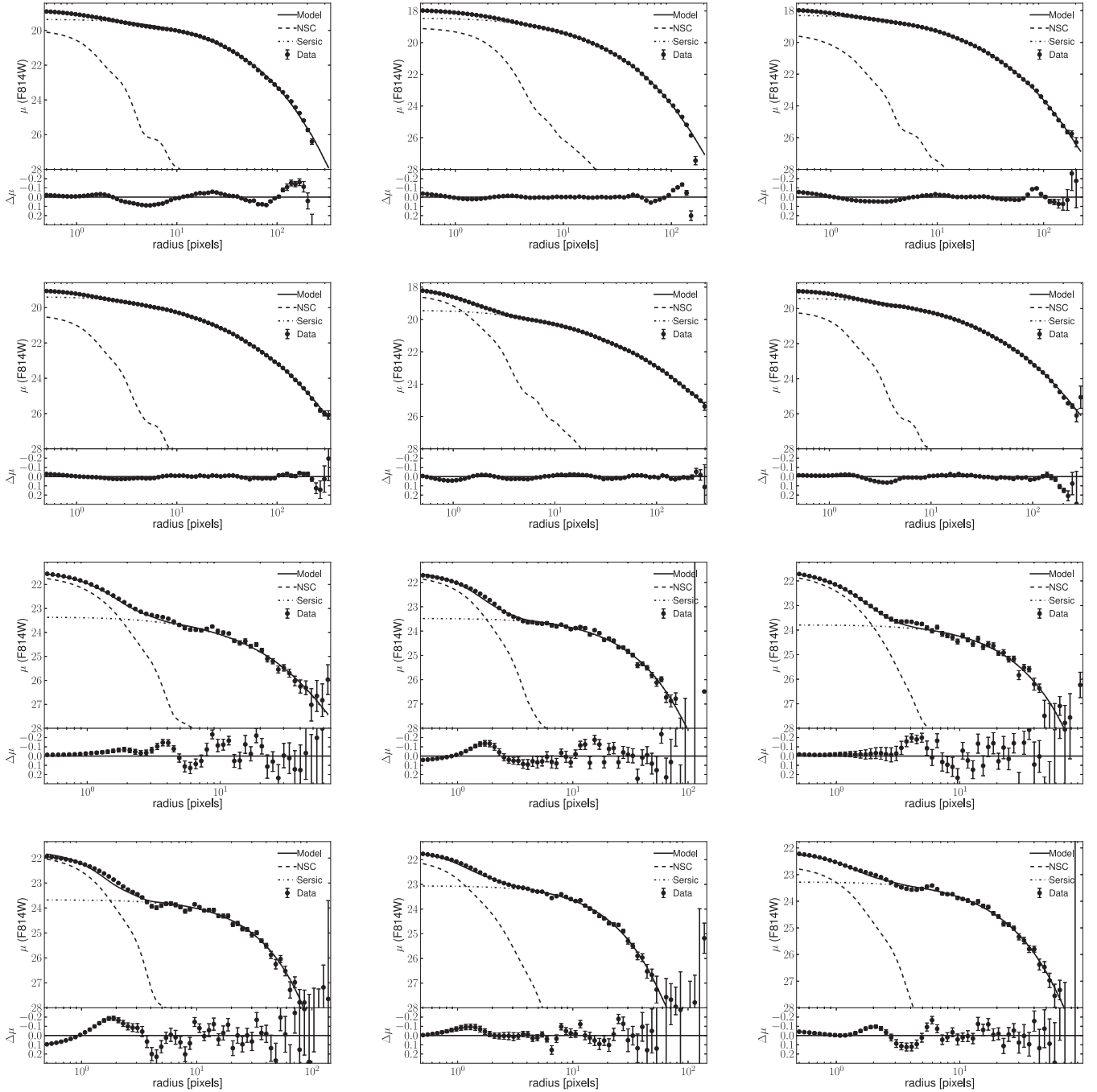


**Table A2.** Sources excluded from the analysis.

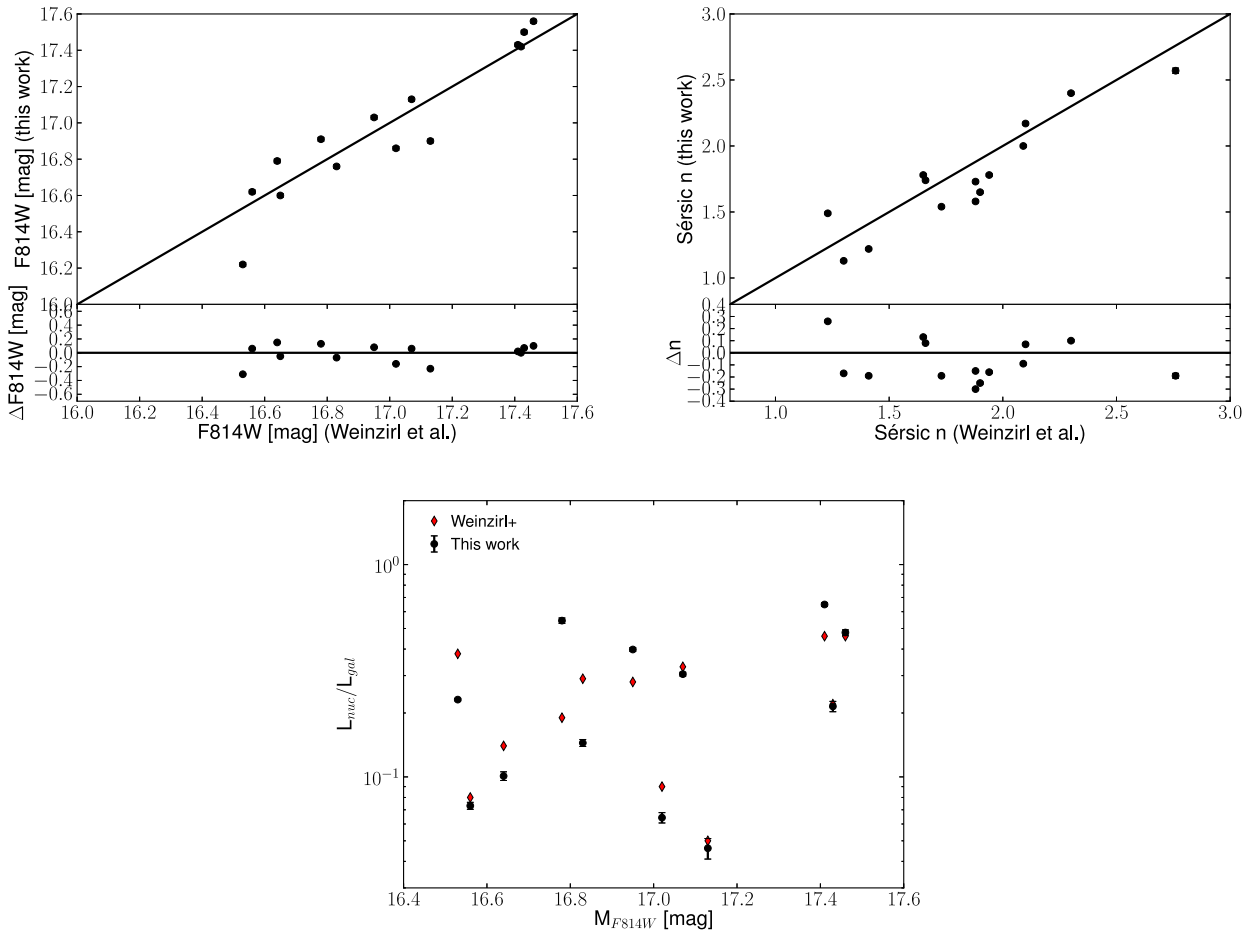
Name	$P_m$	Type	Tile	$z$	RA	Dec.	F814W	Notes
SDSSJ130013.42–280311.8	0	S0	10	8124.0	195.0559	28.0533	17.3	Bright and completely edge-on
SDSSJ130035.42–275633.9	0	S0	22	6925.0	195.1458	27.9428	17.3	This galaxy is not axisymmetric: either barred or two spiral arms
SDSSJ125937.00–280106.9	0	E	12	7195.0	194.9042	28.0186	17.7	The fit is degenerate – centre can be fit in many ways.
SDSSJ125815.27–272752.9	0	S0	45	7615.0	194.5636	27.4647	17.8	Edge-on galaxy
SDSSJ130009.46–275456.3	3	E	24	–	195.0398	27.9156	17.9	Close to bright star
SDSSJ130039.10–280035.5	0	dE	8	5785.0	195.1627	28.0099	18.1	The isophotes of this object are not elliptical. Not clear if central object is NSC
SDSSJ130017.64–275915.1	0	E	16	5966.0	195.0735	27.9876	18.1	Fit requires two outer components + central Sérsic.
CcGV18	0	E	18	6535.0	194.9996	27.9894	20.1	Compact
CcGV9b	0	E	9	6425.0	195.1137	28.0092	19.2	Compact
CcGV1	0	E	1	6775.0	195.1986	28.0927	19.3	Compact
CcGV19b	0	E	19	7075.0	194.9133	27.9985	19.9	Compact
CcGV18	0	E	18	6535.0	194.9996	27.9894	20.1	Compact
CcGV12	0	E	12	7721.0	194.9263	28.0153	20.5	Compact
RB110	0	dE,N	8	7615.0	195.1607	28.016	18.3	Nucleated. In halo of neighbouring galaxy
–	0	dE,N	63	9054.0	194.1588	27.2178	18.6	Galaxy shows spiral arms.
–	1	dE,N	10	–	195.0238	28.0259	18.8	In halo of huge galaxy.
SDSSJ130023.47–280301.9	0	dE	9	6925.0	195.0977	28.0505	19.8	Strong gradient due to neighbouring galaxy
–	1	VLSB	9	–	195.134	28.0376	19.9	Even though this was identified as a galaxy, it looks more like a halo of another galaxy
–	2	dE,N	25	–	194.9245	27.9262	20.1	In halo/edge. Nucleated
–	2	dE	19	–	194.9022	27.9608	20.1	Inside halo
–	2	dE,N	22	–	195.1479	27.9438	20.2	In halo of another galaxy. Nucleation uncertain
–	2	dE,N	33	–	194.8487	27.8451	20.3	Galaxy is in edge – fit looks decent but nucleation is uncertain.
SDSSJ125845.91–274655.5	2	dE,N	75	–	194.6912	27.7823	20.3	Galaxy is edge-on/irregular
SDSSJ125942.92–275954.6	0	dE,N	19	8274.0	194.9288	27.9984	20.4	Strong gradient in the sky due to nearby galaxy/Irregular
–	2	dE	25	–	194.9535	27.9155	20.4	In area of low S/N
–	0	S0	18	8814.0	194.9988	27.9983	20.5	Either a bar or offset NSC
–	1	VLSB	19	–	194.9158	27.9905	20.5	VLSB – not nucleated
–	3	dE	23	–	195.0918	27.8987	20.6	In halo
SDSSJ130039.76–280601.9	1	VLSB	2	–	195.1658	28.1008	20.9	Nucleated but difficult to fit.
SDSSJ125850.42–274445.7	2	dE,N	75	–	194.71	27.746	20.9	The fit is dubious – asymmetric residuals around the NSC.
–	2	dE	15	–	195.1595	28.0024	21.0	Covered by a star
–	3	dE,N	25	–	194.9599	27.9216	21.0	Crowded area: difficult to obtain good fit
SDSSJ130011.81–280504.0	2	dE,N	3	–	195.049	28.0846	21.1	Blended with neighbour. Likely nucleated.
–	2	dE	23	–	195.0737	27.9368	21.1	In halo
SDSSJ126944.76–275807.1	0	dE	18,19	9623.0	194.9363	27.9685	21.2	VLSB in edge of frame
–	3	dE,N	13	–	194.8752	28.0439	21.6	Extremely elongated – difficult to fit and probably background
SDSSJ130000.97–275929.5	1	dE,N	18	–	195.0042	27.9929	21.6	In halo of elliptical galaxy – nucleated
–	3	dE,N	13	–	194.8776	28.0223	21.7	In halo – nucleated
SDSSJ130037.83–275840.9	0	dE	15	4684.0	195.1576	27.9779	21.7	Compact source
–	2	dE	33	–	194.9011	27.8853	21.7	VLSB – nonnucleated
–	2	dE,N	19	–	194.9108	27.9496	21.8	Unfittable – in halo (nucleated)
–	2	dE	19	–	194.8825	28.0015	22.0	Irregular galaxy
SDSSJ125832.93–272406.5	3	dE,N	46	–	194.637	27.402	22.0	Irregular galaxy probably background
–	2	VLSB	19	–	194.9204	27.9548	22.1	Very crowded area – nucleation and centre uncertain.
–	3	dE	22	–	195.1553	27.9236	22.1	In the spokes of a bright foreground star
–	3	dE	2	–	195.1355	28.0606	22.2	Fit affected by nearby bright galaxy
–	3	dE,N	12	–	194.9351	28.0413	22.2	Poor data quality – possible NSC is masked out as cosmic.
–	2	dE,N	19	–	194.9068	27.9686	22.2	In halo
–	3	dE,N	10	–	195.0617	28.0085	22.4	Centre uncertain
–	3	dE,N	24	–	195.0102	27.9497	22.4	This galaxy is misclassified as dE,N

Table A2 – continued

Name	$P_m$	Type	Tile	$z$	RA	Dec.	F814W	Notes
–	3	dE	24	–	195.0385	27.9076	22.4	Difficult to fit
–	2	dE,N	45	–	194.6363	27.4596	22.4	In edge and halo
–	2	dE	10	–	195.0524	28.0085	22.5	Difficult to fit
–	2	dE,N	19	–	194.9121	27.9789	22.5	Very difficult to fit, in 4874 halo
–	3	dE	22	–	195.1334	27.9519	22.5	Difficult to fit



**Figure A1.** Profiles of the brightest and faintest galaxies in the sample. From left to right, top to bottom: SDSSJ130018.54–280549.7, LEDA126789, SDSSJ130041.19–280242.4, SDSSJ130034.42–275604.9, RB068, LEDA126815, COMAi125932.883p275800.05, COMAi125929.995p275348.12, COMAi125952.543p275824.21, COMAi13021.712p275650.16, COMAi125944.017p275615.29, COMAi125924.938p275320.35



**Figure C1.** Comparison of the structural parameters for galaxies in common with Weinzirl et al. (2014). Upper-left panel:  $F814W$  host galaxy magnitude comparison. Upper-right panel: Sérsic index from this work plotted against the Sérsic index found by Weinzirl et al. Lower panel: fraction of the total luminosity emitted by the nucleus as a function of host galaxy magnitude.

modules of the FFTW<sup>34</sup> library. We note that both the model and surrounding zero-pads are sufficiently big. After convolution, an optional sky background is added to the model.

### B3 Comparison with data

For each pixel, we calculate the probability that it was generated by the model. For this, we assume Gaussian errors. The total likelihood is the product of the likelihoods of the individual pixels, except those that are flagged as bad.

### B4 Likelihood exploration

The likelihood is explored using the nested sampling algorithm (Skilling 2004). The benefit of using nested sampling over ordinary MCMC codes is its ability to accurately calculate the Bayesian evidence for a model. The algorithm carries out a one-dimensional integral of likelihood samples over prior mass. New likelihood samples are generated using the ellipsoidal sampling algorithm of (Mukherjee, Parkinson & Liddle 2006). We do modify the standard way of calculating the mean and variance of the live sample for angle variables (in our case the PA of the galaxy). Instead, we calculate the mean angle of a collection of angles  $\theta_i$  by the arctangent

of the mean cosine and sine of the angles:

$$\langle \theta \rangle = \tan^{-1} \left( \frac{\langle \sin(\theta_i) \rangle}{\langle \cos(\theta_i) \rangle} \right), \quad (\text{B1})$$

where  $\langle \dots \rangle$  denotes an average.

After the nested sampling has finished, we calculate the mean and variance of each free parameter by taking the mean and second moment of the posterior weight, as advertised in Skilling (2004). The code automatically generates a model and residual image.

## APPENDIX C: COMPARISON WITH WEINZIRL ET AL.

In Fig. C1, we compare the structural parameters fitted by Weinzirl et al. (2014) with GALFIT for the 15 galaxies in common. We note that we fit the NSC in some cases by a Gaussian and use a different PSF, which leads to a different luminosity than the point source (used in all cases) by Weinzirl et al. For the brightest source in common, we find a difference in the host galaxy magnitude of 0.3 mag. The average difference in magnitude is  $0.01 \pm 0.03$  mag, with a standard deviation of 0.13. Excluding the 0.3 mag offset source lowers the standard deviation to 0.1 mag. The difference in the logarithm of the Sérsic index  $\ln(n)$  is  $0.04 \pm 0.03$ , with a standard deviation of 0.10. Despite the use of different codes to fit

<sup>34</sup> <http://www.fftw.org>

the galaxies, different masking schemes, different weighting and different PSFs, the comparison is excellent.

## SUPPORTING INFORMATION

Additional Supporting Information may be found in the online version of this article:

**Table A1.** Structural parameters of galaxies used in the paper. (<http://mnras.oxfordjournals.org/lookup/suppl/doi:10.1093/mnras/stu1906/-/DC1>).

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